

Bilayer excitons in strongly correlated materials

Louk Rademaker, Kai Wu, Jan Zaanen and Hans Hilgenkamp

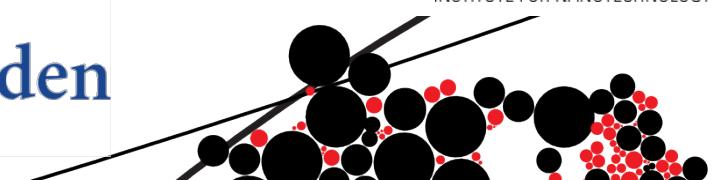
rademaker@lorentz.leidenuniv.nl

Institute-Lorentz, University of Leiden and MESA+ institute for Nanotechnology, University of Twente

Tallahassee – FSU/NHMFL – 1 February 2013



UNIVERSITY OF TWENTE.
Universiteit Leiden



Introduction: Exciton condensation

Exciton:

Bound electron-hole pair

Bose condensation

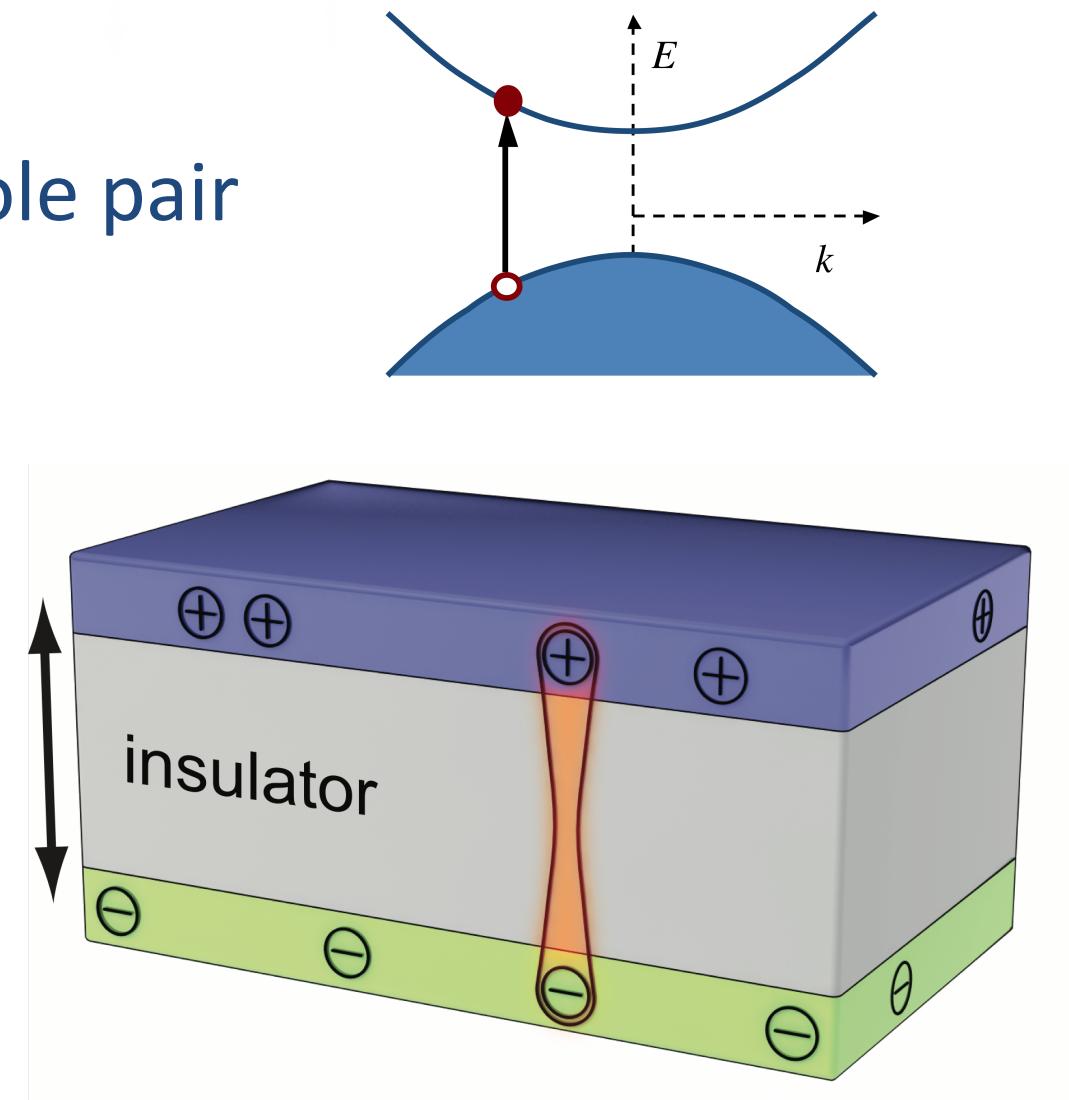
(BCS-style)

Good: Coulomb

Bad: Annihilation

Solution:

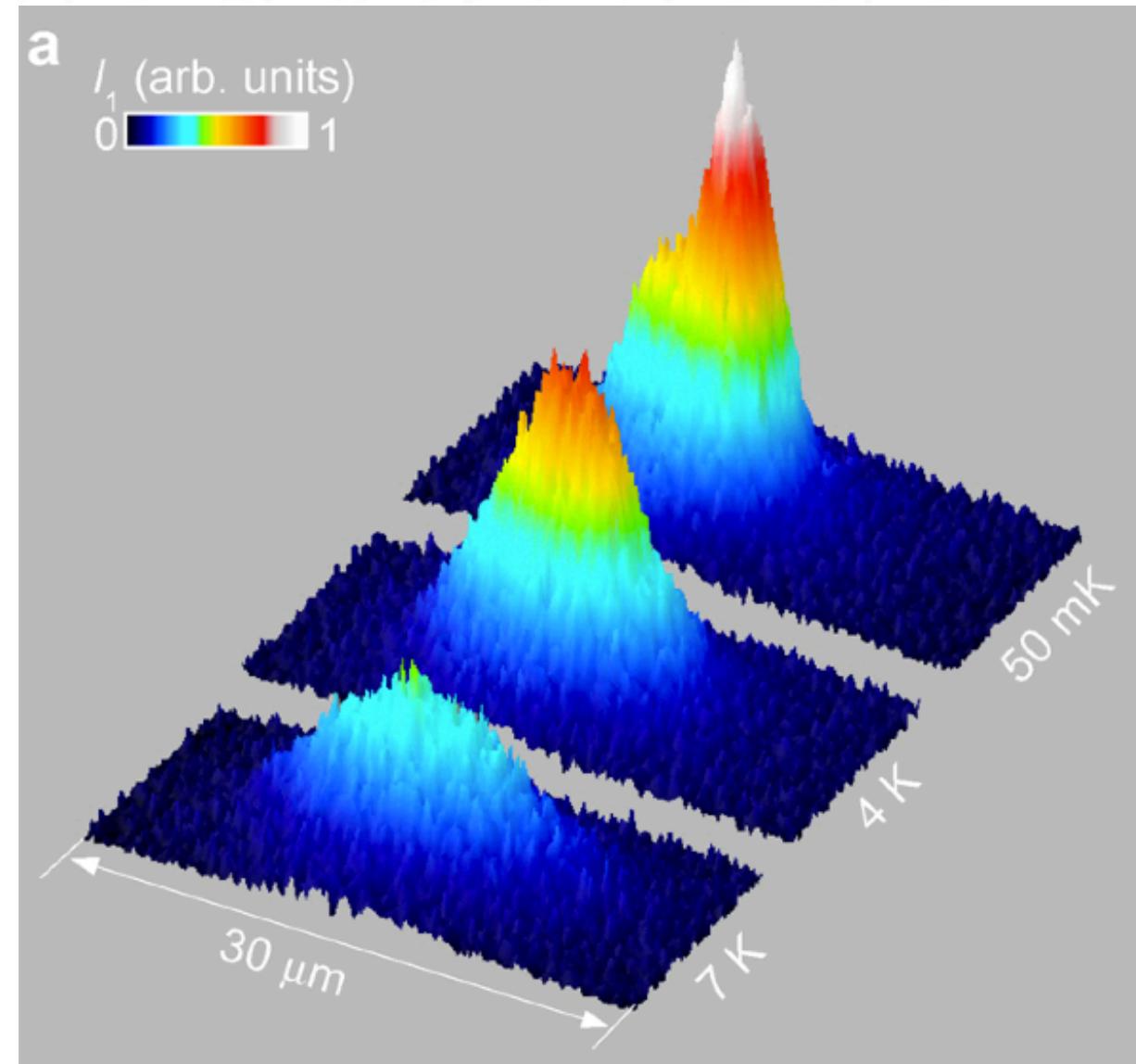
Bilayer systems



Introduction: Exciton condensation

Pumped
bilayer
quantum
wells

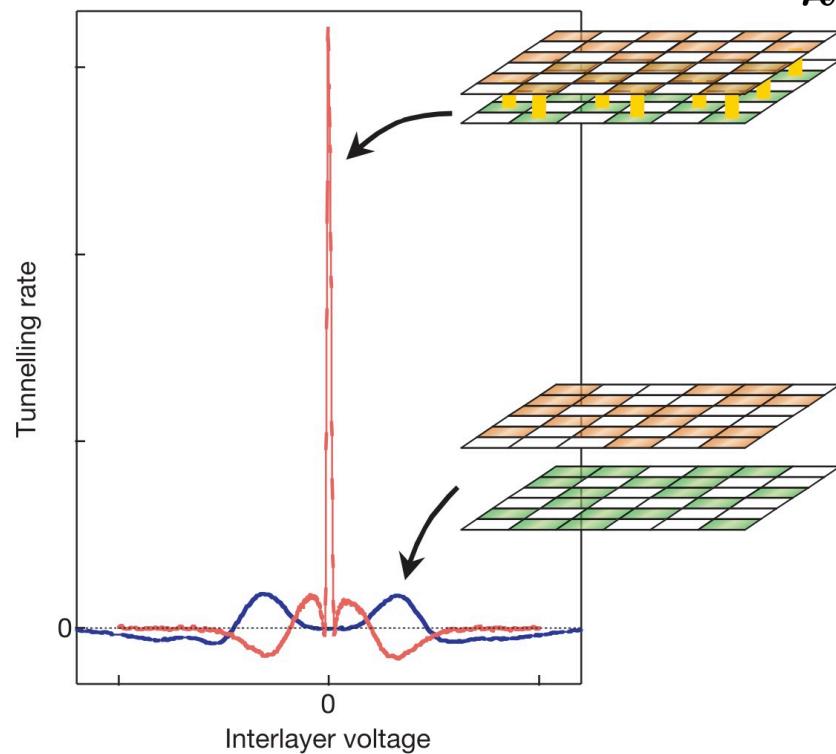
Ref: High et al, Nano Lett 2012



Introduction: Exciton condensation

Directly measure order parameter

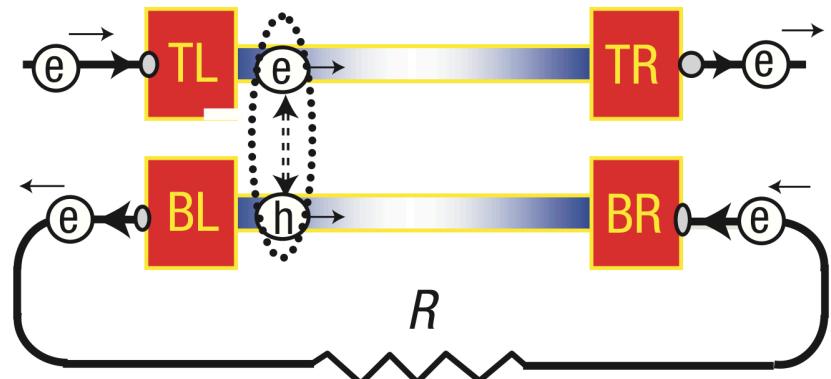
$$\Delta = \langle c_k^\dagger c_2^\dagger \rangle_k \rangle$$



Ref: Eisenstein & MacDonald, Nature 2004

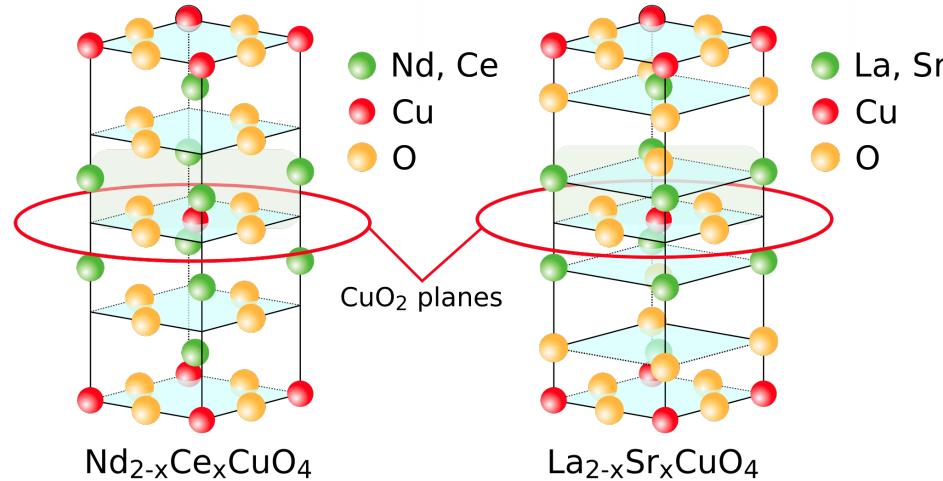
Counterflow superfluidity

Drag-counterflow (D-CF)



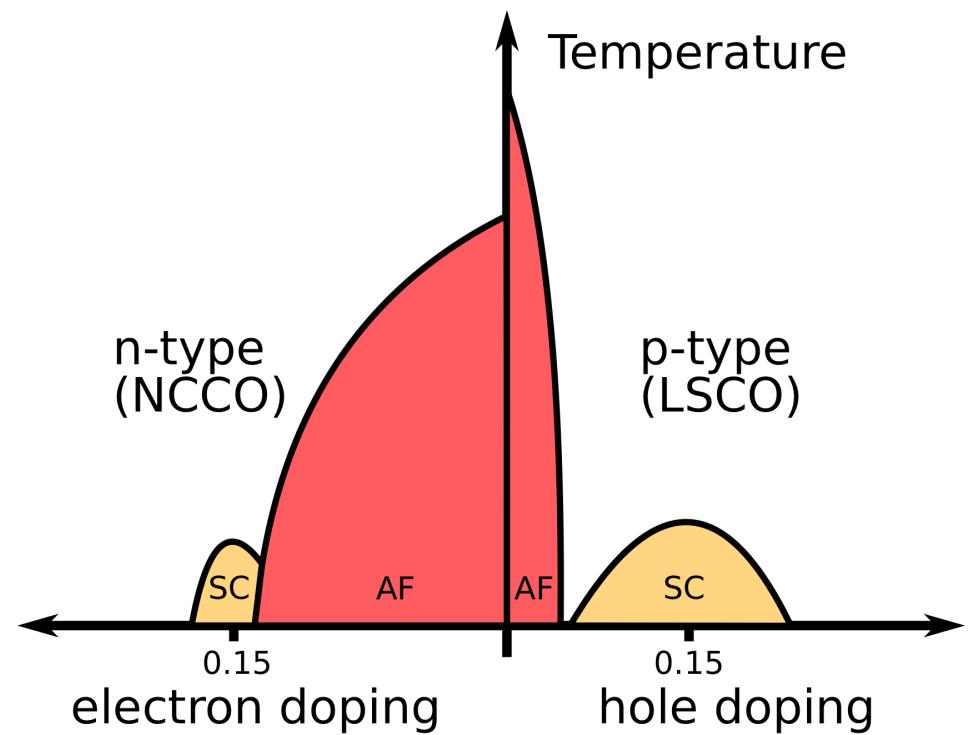
Ref: Su & MacDonald, Nat Phys 2008

Introduction: Strongly correlated materials



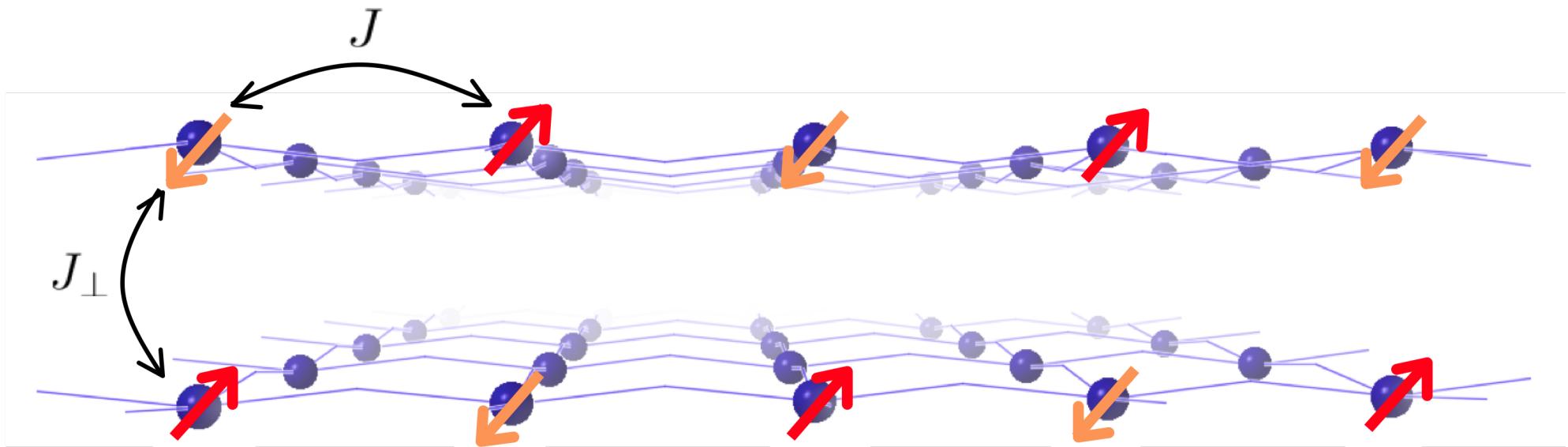
Example: **cuprates**
Quasi 2d CuO₂ planes

Mott insulator:
local moments
failure of band theory



Introduction: Strongly correlated materials

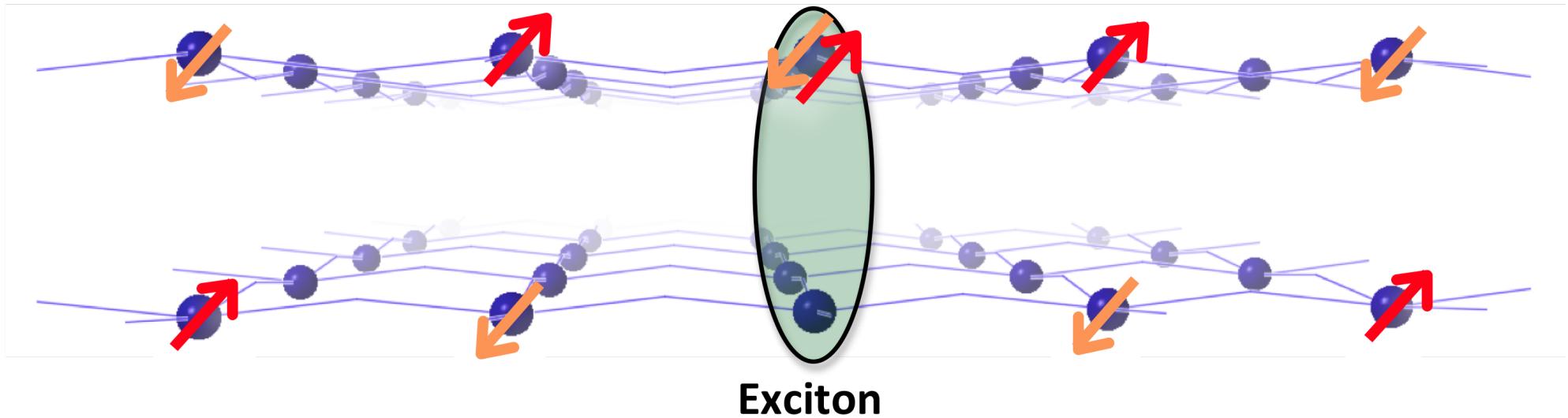
Antiferromagnetic order at half-filling



Heisenberg interactions $J \vec{S}_i \cdot \vec{S}_j$

Excitons in strongly correlated materials

Excitons are: *doublon-holon pairs*



Condensation? Other phenomena?

Work with Jan Zaanen



& Hans Hilgenkamp



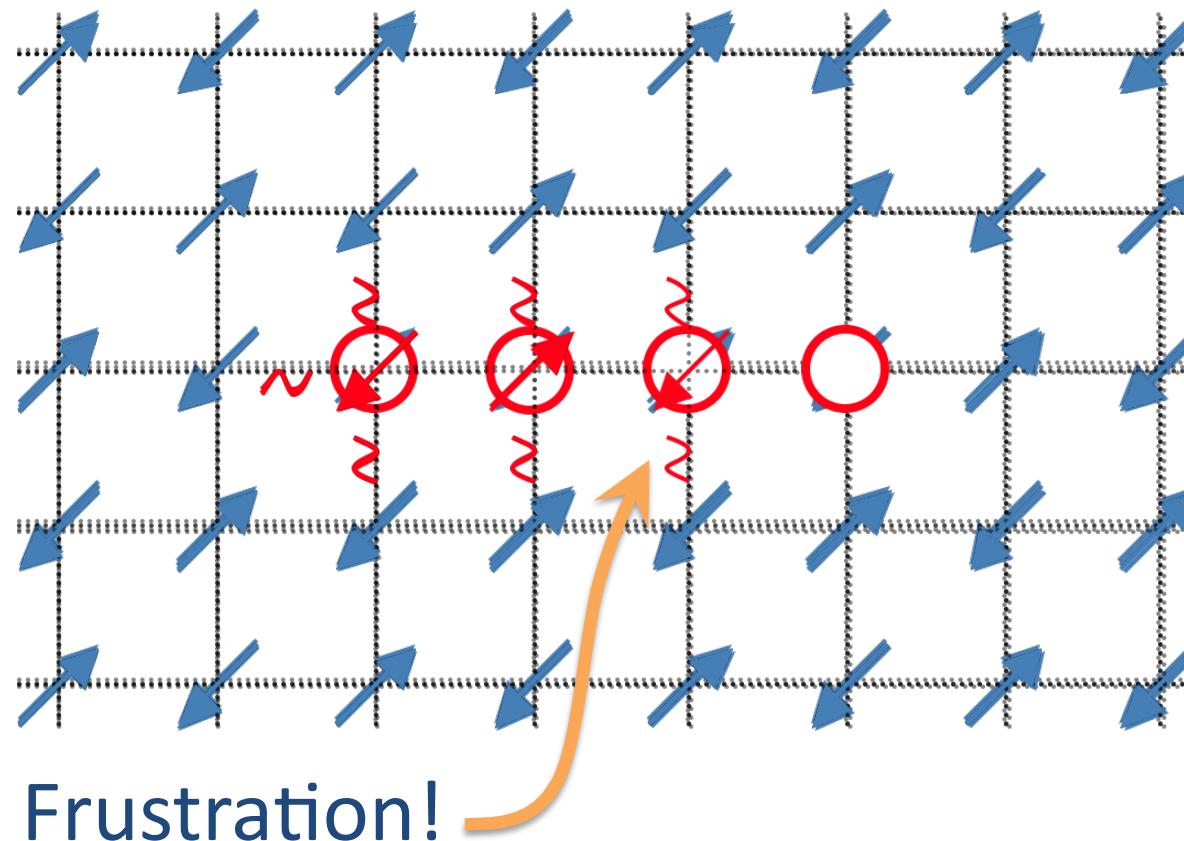
Lay-out of the rest of this talk

Two main questions:

1. What are the **dynamics** of a **single exciton**?
2. What is the **phase diagram** with a finite **density** of excitons?

Dopant in Mott insulator: Frustration

Moving hole in a Mott insulator:



Ising confinement

Linear potential: Ladder spectrum!

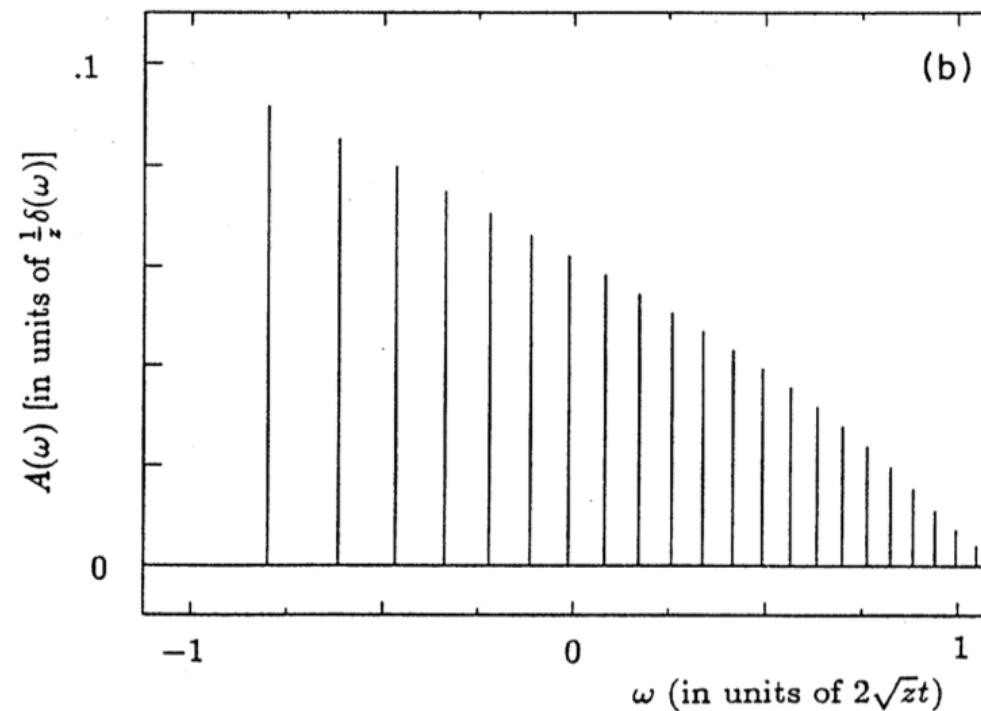


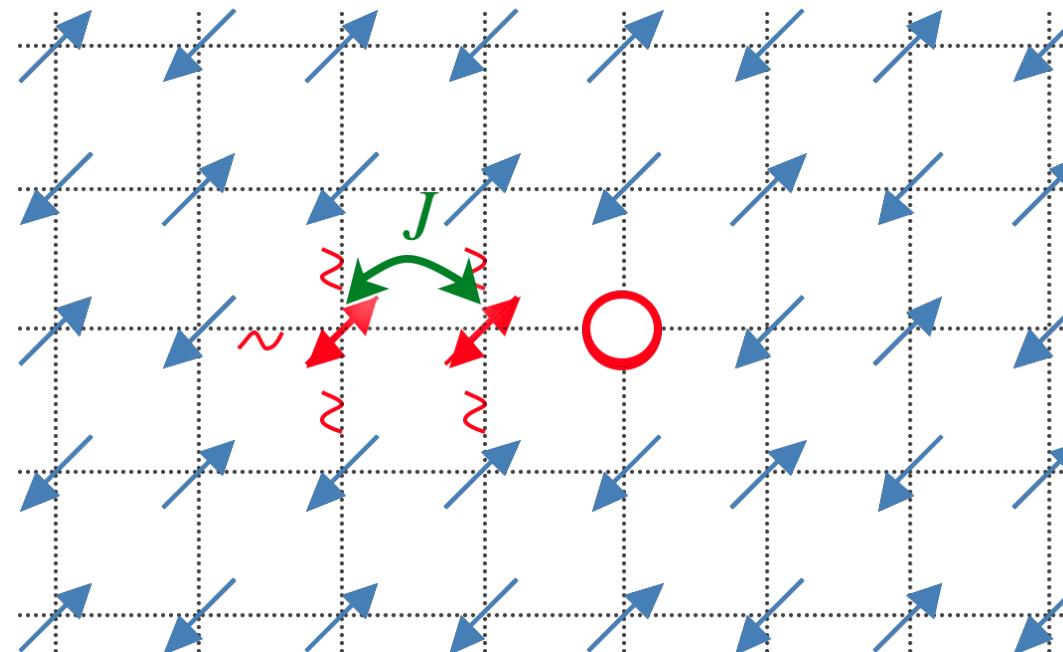
FIG. 3. Hole spectral function in the $J_{\perp}=0$ limit. (a) The limit $J_z=0$. (b) $J_z/t=0.1$. The vertical lines represent delta functions with weight specified by their height.

Ref: Kane, Lee & Read, PRB 1989

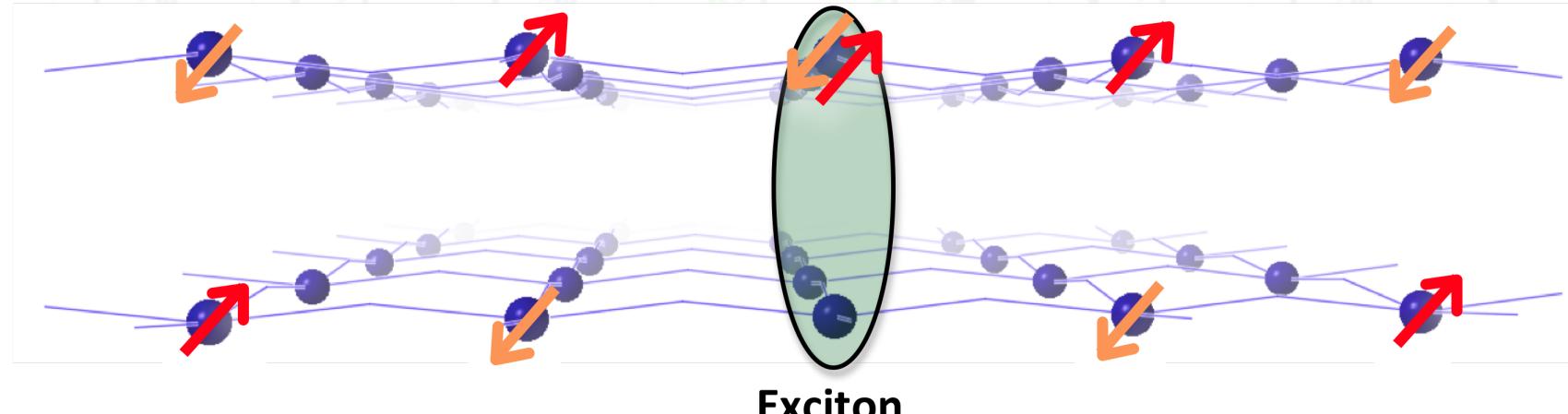
Quantum corrections: no confinement

Heisenberg antiferromagnet

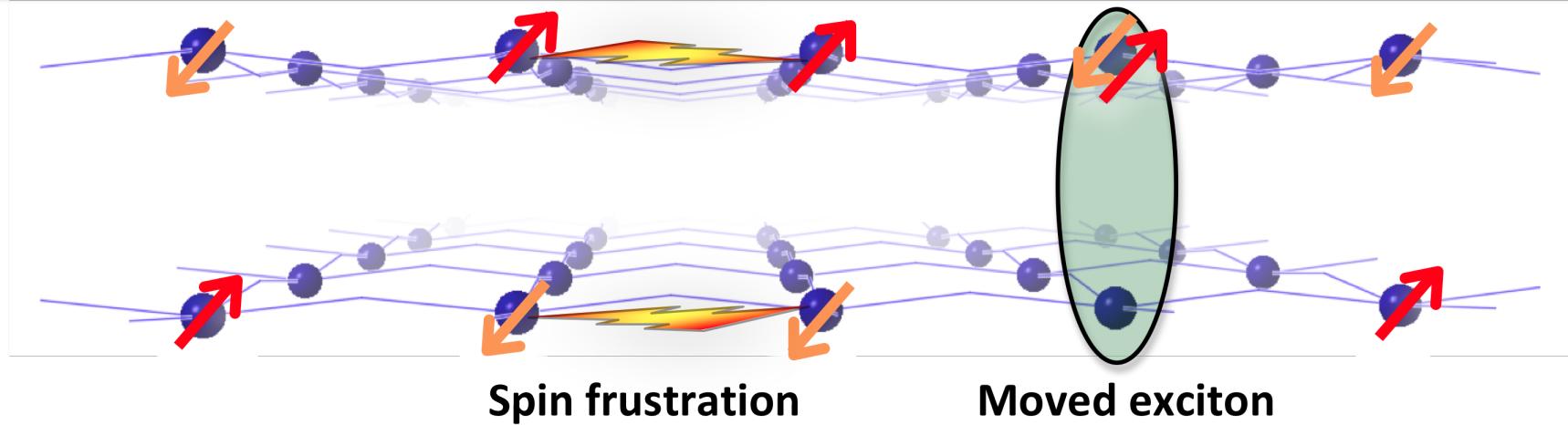
Frustration is repaired by $JS_i^+S_j^-$



Back to the bilayer: *exciton t-J model*



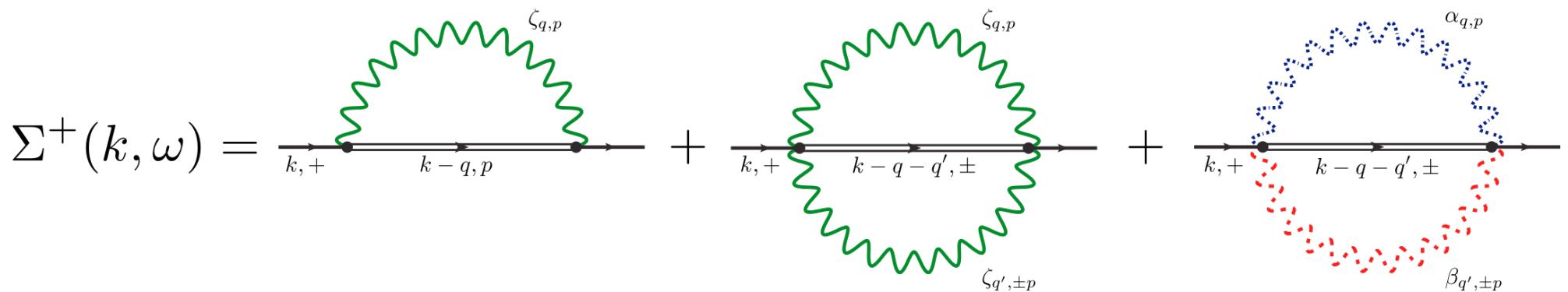
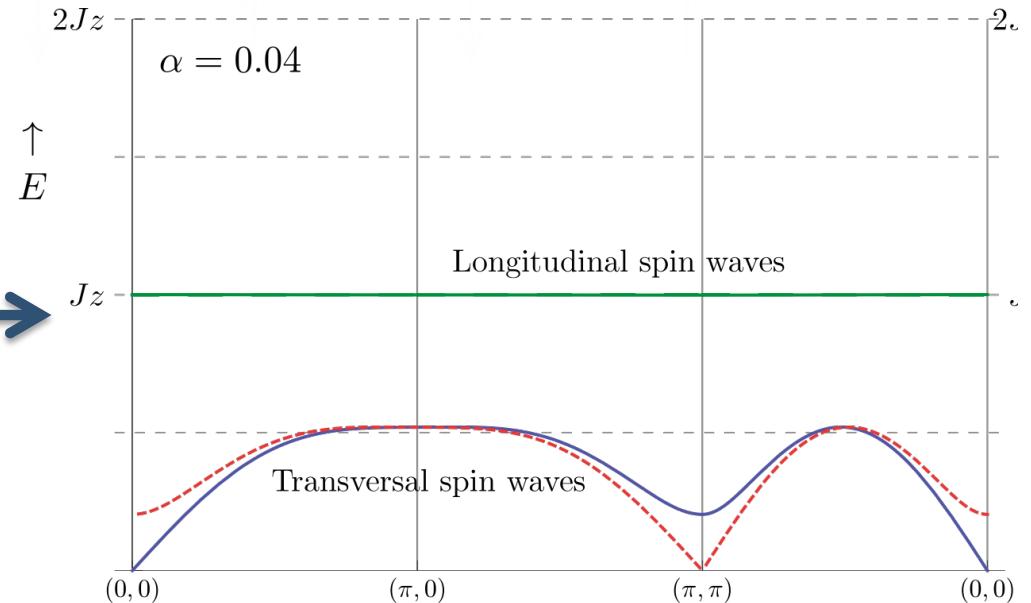
$$H = J \sum_{\langle ij \rangle l} \mathbf{S}_{il} \cdot \mathbf{S}_{jl} + J_\perp \sum_i \mathbf{S}_{i1} \cdot \mathbf{S}_{i2} - t \sum_{\langle ij \rangle} |E_j\rangle \left(|0\ 0\rangle_i \langle 0\ 0|_j + \sum_m |1\ m\rangle_i \langle 1\ m|_j \right) \langle E_i|$$



Spin wave theory

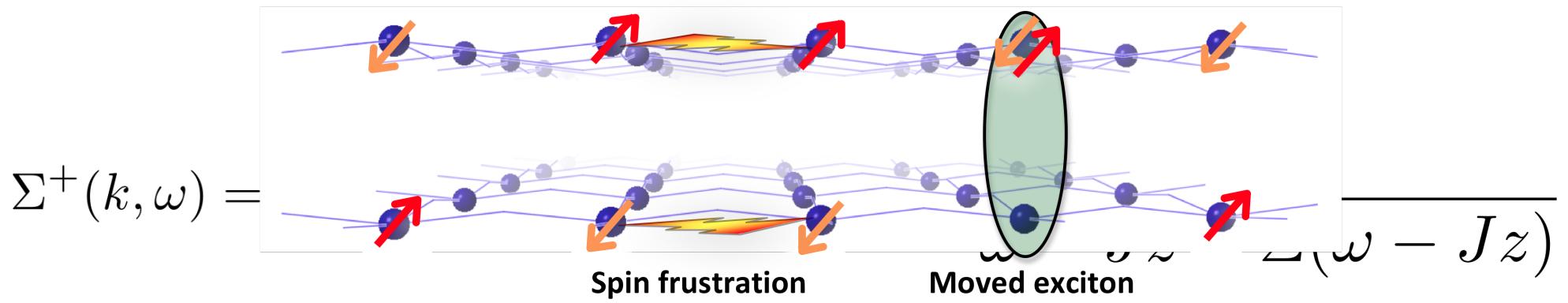
Moving exciton
excites and absorbs
spin waves →

Exciton self-energy:



Dynamical frustration once again

Have to excite the **gapped spin wave** to move

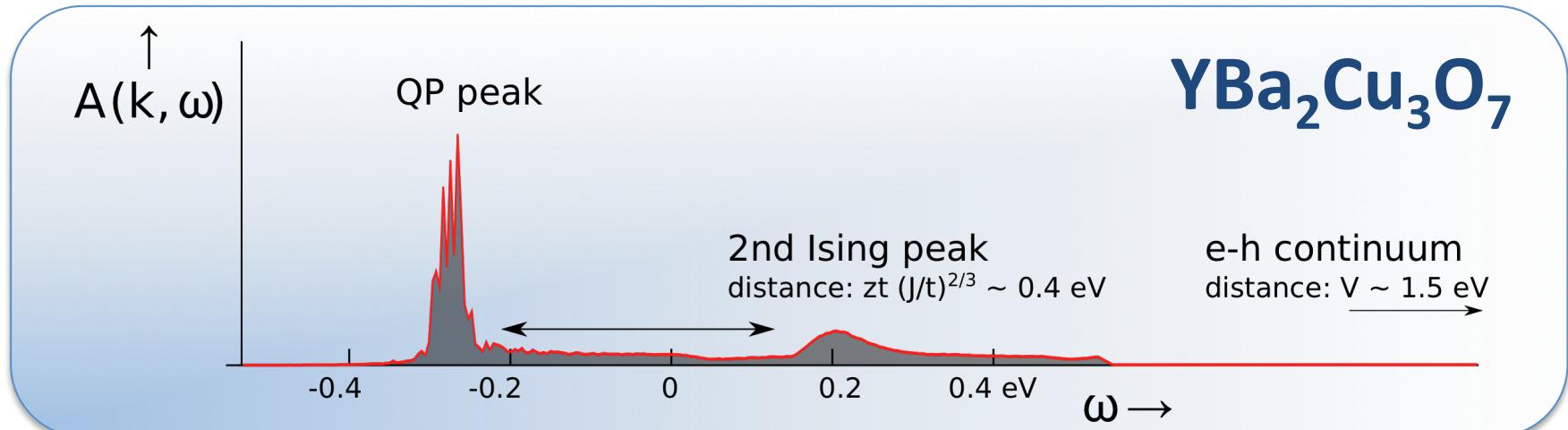
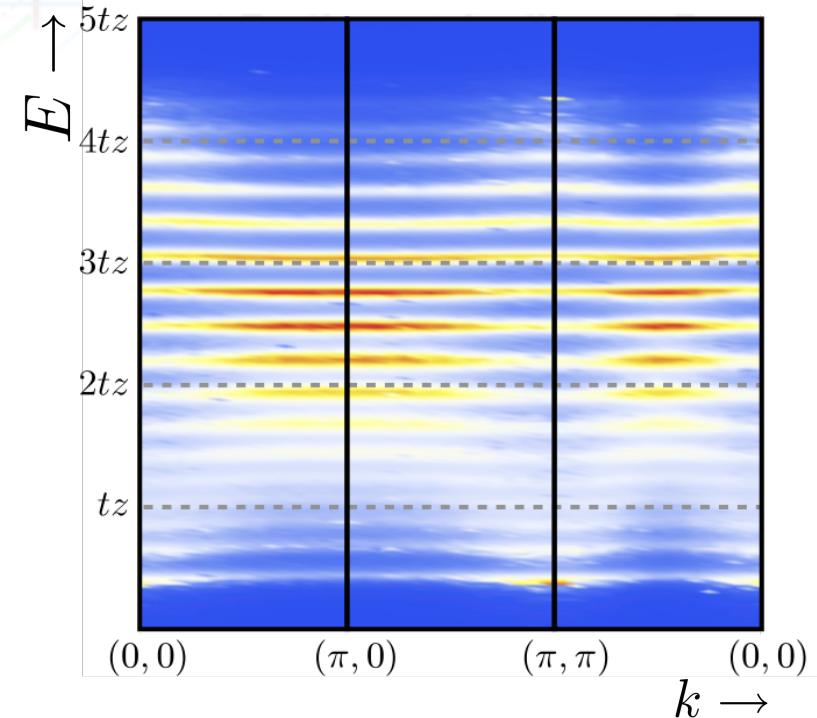
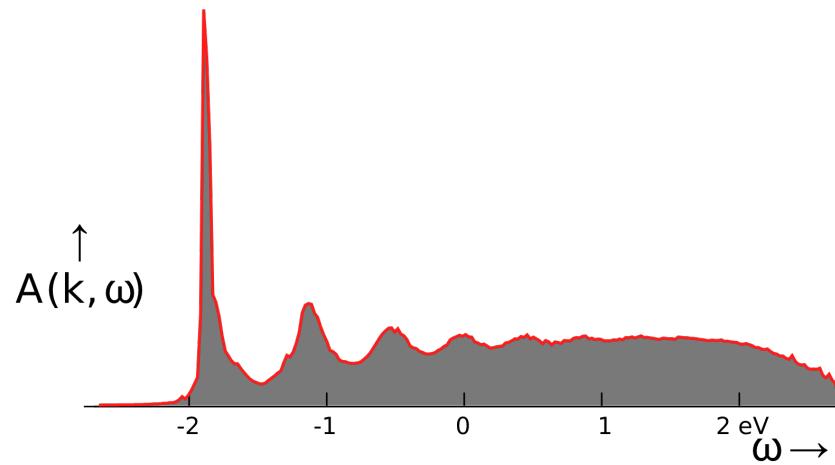


Same self-energy as Ising confinement!

Observable (at $k=0$) in optical absorption

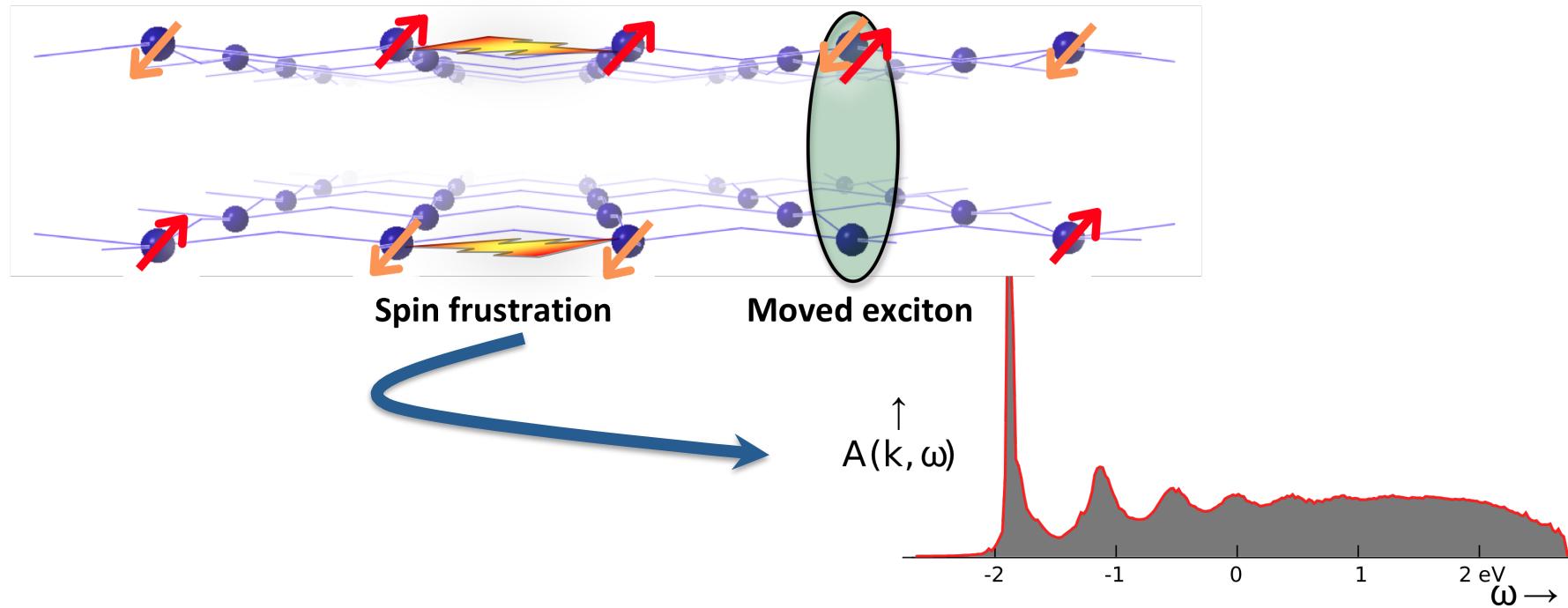
Results: Exciton spectral function

Ladder spectrum:



Summarizing single exciton results

Enhanced dynamical frustration



Ref: Rademaker, Wu, Zaanen & Hilgenkamp, EPL 2012;
Rademaker, Wu, Zaanen, NJP 2012

Hardcore bosons on a lattice

First: forget about the spin background

Do hardcore bosons on a lattice $E^+ = |1\rangle\langle 0|$

$$H = -t \sum_{\langle ij \rangle} (E_i^x E_j^x + E_i^y E_j^y) - \mu \sum_i E_i^z + V \sum_{\langle ij \rangle} E_i^z E_j^z$$

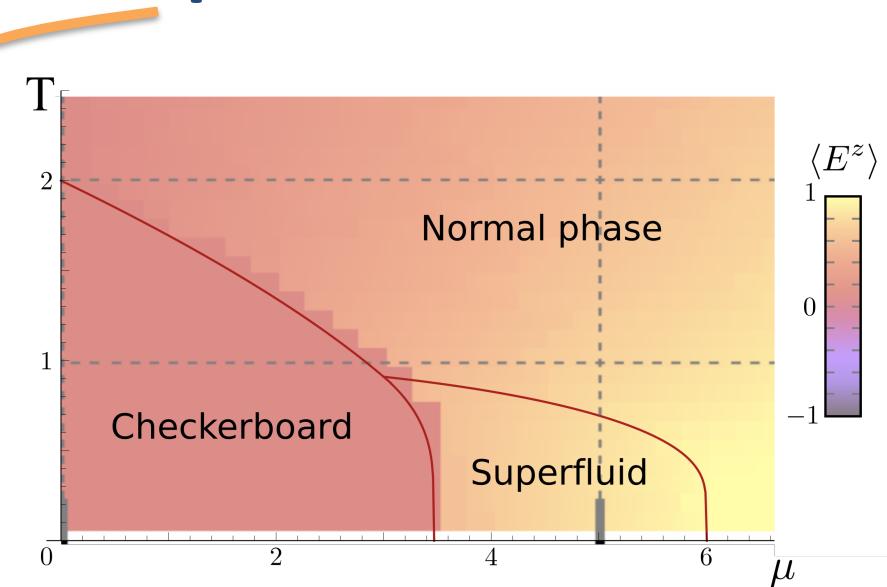
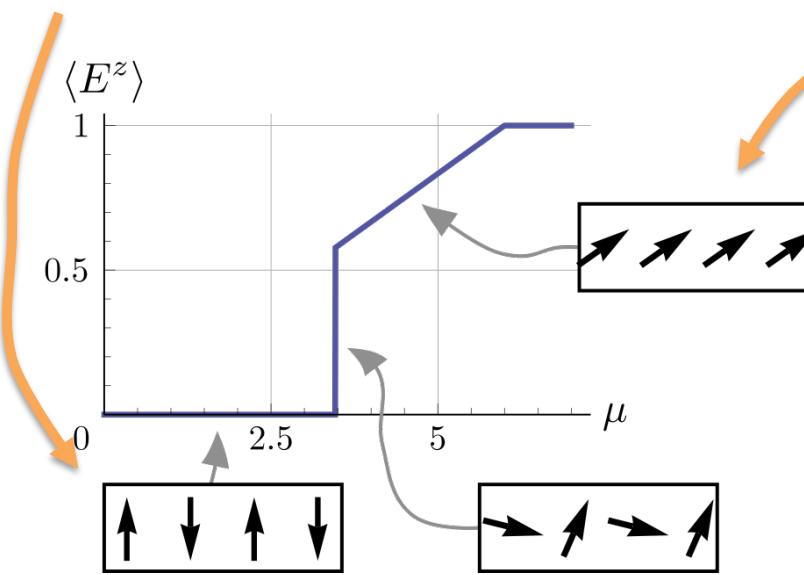
Ground state: variational wavefunction

$$|\Psi\rangle = \prod_i (\cos \theta_i e^{i\psi_i} |1\rangle_i + \sin \theta_i |0\rangle_i)$$

Hardcore bosons on a lattice

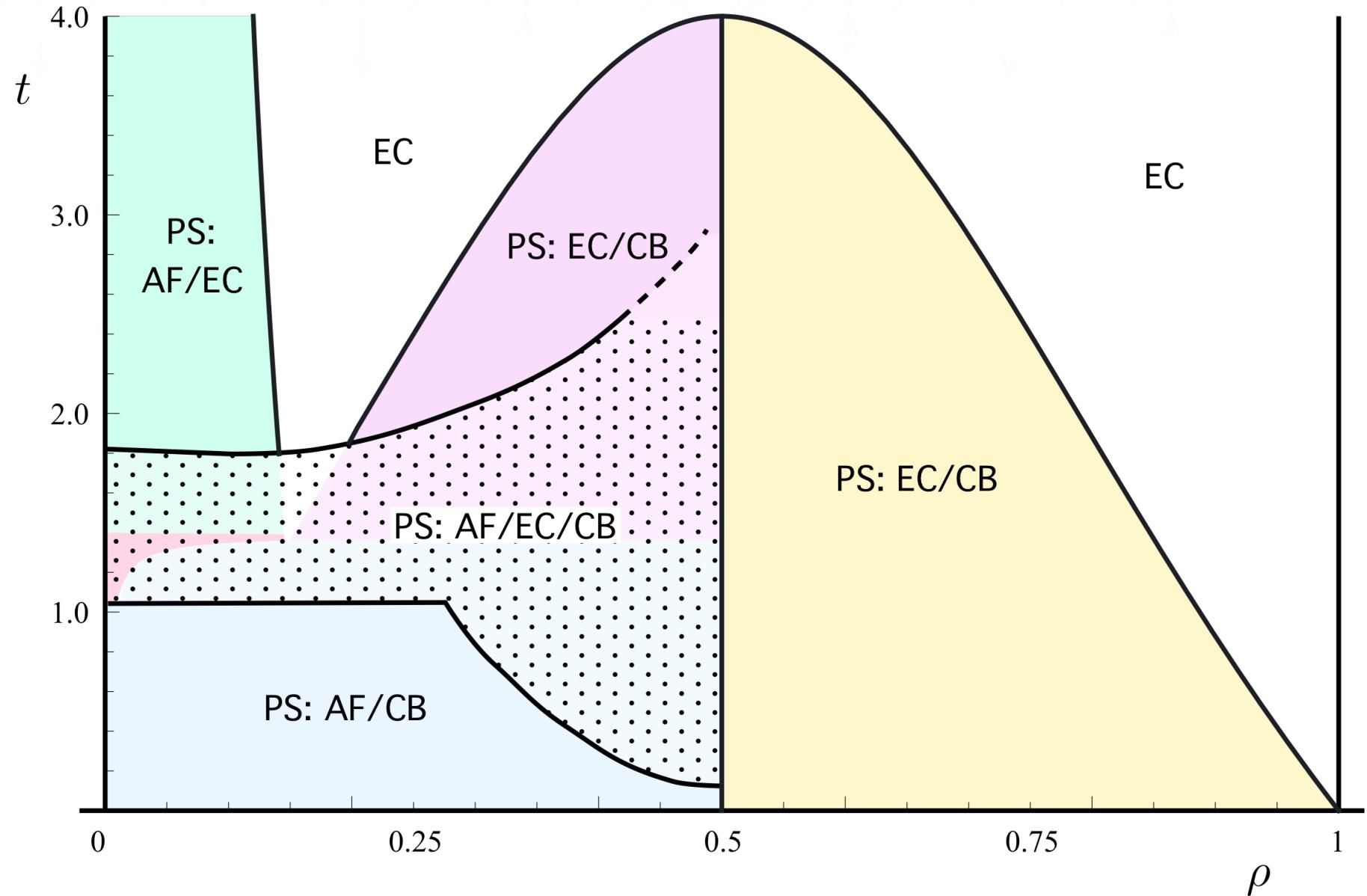
Two phases:

Checkerboard / solid and Superfluid



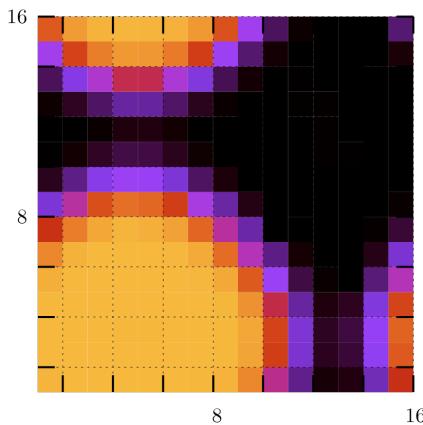
Separated by 1st order transition ('spin-flop')

Ground state phase diagram

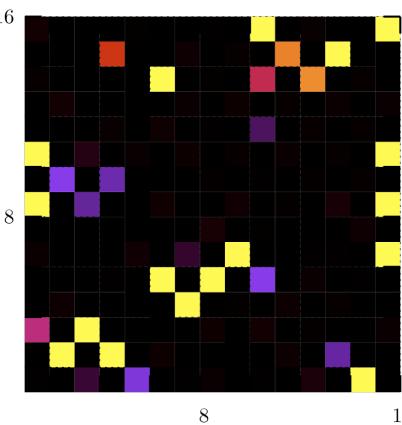


Phase separation: Monte Carlo results

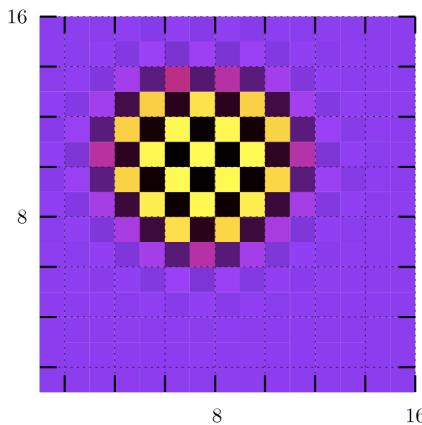
a. $\rho = 0.05, t = 2.3 \text{ eV}$



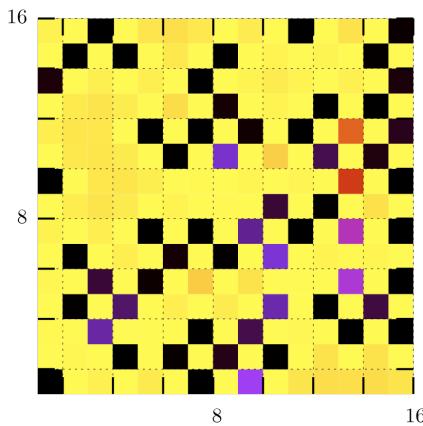
b. $\rho = 0.1, t = 0.1 \text{ eV}$



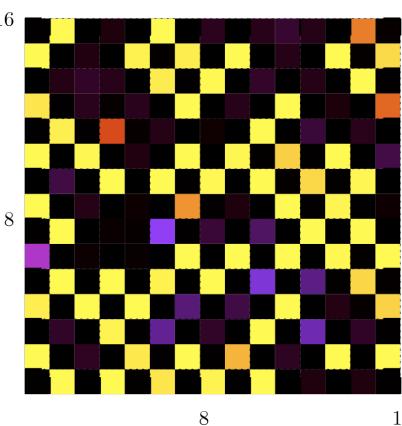
c. $\rho = 0.25, t = 2.3 \text{ eV}$



d. $\rho = 0.75, t = 0.5 \text{ eV}$



e. $\rho = 0.3, t = 1.5 \text{ eV}$



Exciton density scale:

1.0 (0.14 for a.)

0.5 (0.07 for a.)

0.0

The exciton superfluid

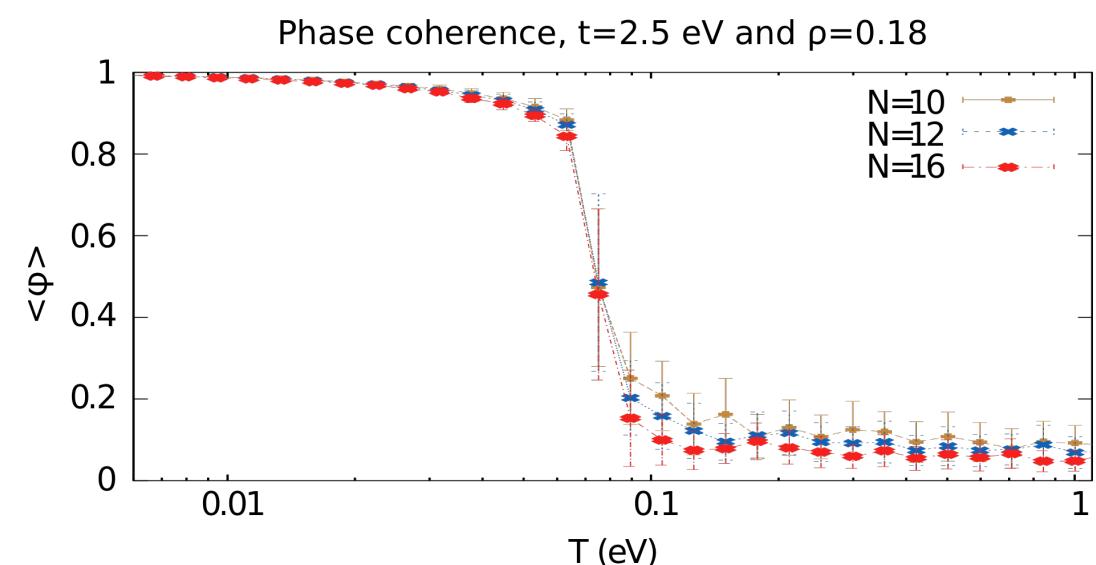
Ground state wavefunction

$$|\Psi\rangle = \prod_i \left(\sqrt{\rho}|E_i\rangle + \sqrt{1-\rho}|0\ 0\rangle_i \right)$$

Anomalous interlayer tunneling

$$\langle \sum_{\sigma} c_{i1\sigma}^{\dagger} c_{i2\sigma} \rangle = \sqrt{2\rho(1-\rho)}$$

Kosterlitz-Thouless
transition:



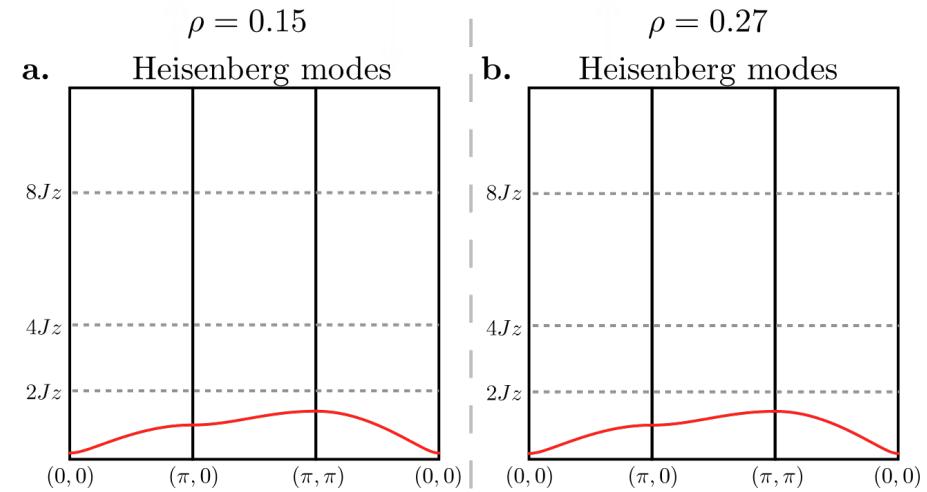
Exciton-spin dynamics in the superfluid



Superfluid = quantum paramagnet

So: propagating triplet excitations

Strong exciton-spin coupling shown in renormalization of triplet bandwidth



Why does triplet kinetic energy increases?

Heisenberg ‘hopping’ of triplets

$$J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j \rightarrow -J \sum_{\langle ij \rangle} t_i^\dagger t_j$$

Exciton hopping

$$\langle e^\dagger \rangle \rightarrow \sqrt{\rho_{SF}}$$

$$-t \sum_{\langle ij \rangle} e_j^\dagger e_i t_i^\dagger t_j$$

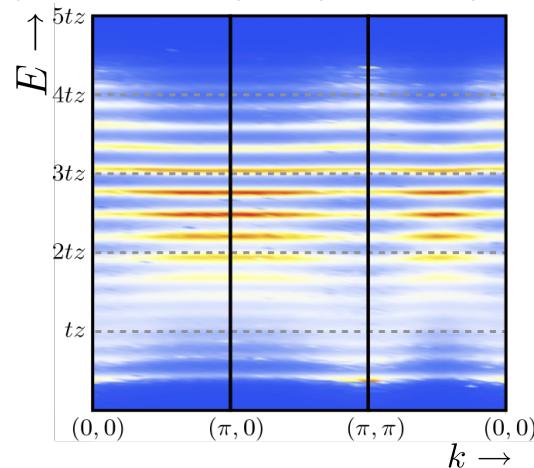
Energy scales: $J \sim \frac{t_e^2}{U}$ $t \sim \frac{t_e^2}{V}$ $U > V \rightarrow t > J$

Thus enhancement...

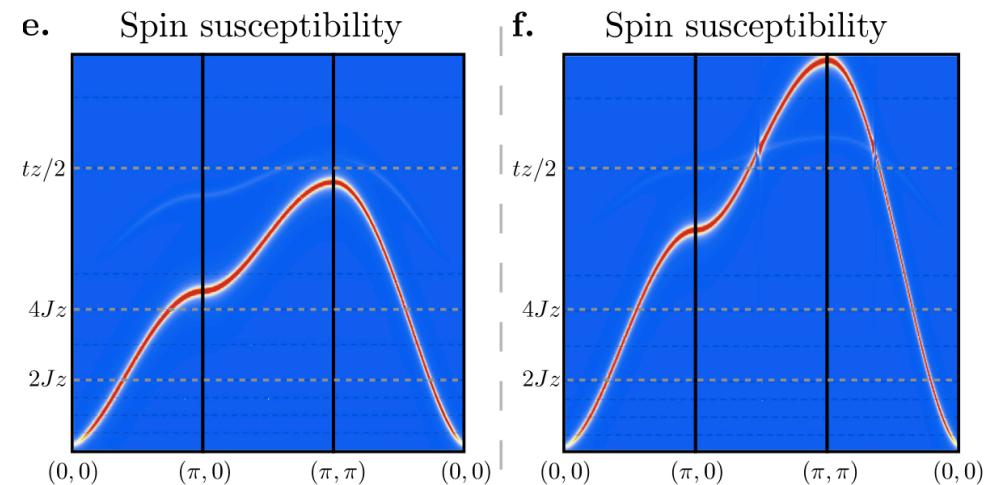
Summary

Magnetic order
prevents motion of
exciton

Ref: Rademaker, Wu, Zaanen & Hilgenkamp,
EPL 2012; Rademaker, Wu, Zaanen, NJP 2012



Exciton condensation
helps motion of
triplets



But: phase separation is all around!

thank you for your attention!