

Exciton condensation in strongly correlated electron bilayers

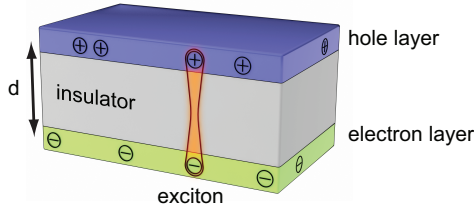
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Our idea: combine bilayer excitons and strongly correlated electron systems

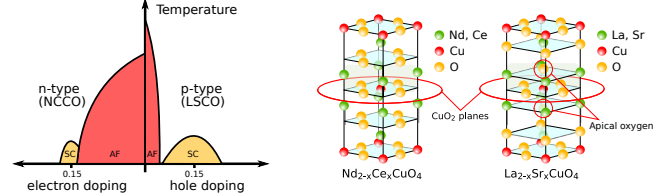
A **bilayer exciton** is an exciton where the hole and electron are spatially separated in two different layers. An insulating layer in between prevents annihilation of the excitons.

Because of their long life-time, bilayer excitons are ideal candidates for realizing a Bose condensate of excitons. [1]



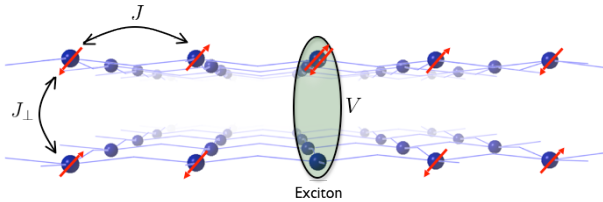
The high-temperature superconducting cuprates are an example of layered **strongly correlated electron systems**.

Because of their quasi-2D nature, strongly correlated electron systems are ideal candidates for bilayer excitons.



What is a bilayer exciton in a strongly correlated background?

In strongly correlated electron systems the electrons are localized, and standard electronic band theory does not apply.



Sideview of a strongly correlated bilayer. In red the localized electron spin is shown, with antiferromagnetic order.

The interaction between the localized electron spins are governed by an antiferromagnetic Heisenberg interaction.

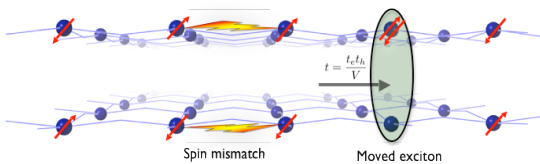
Now a bilayer exciton is the bound state of a doubly occupied and a vacant site.

This leads to the following model Hamiltonian:

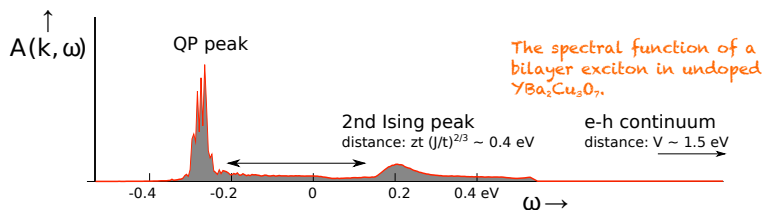
$$\mathcal{H} = J \sum_{\langle ij \rangle} \mathbf{S}_{i1} \cdot \mathbf{S}_{j1} + J_{\perp} \sum_i \mathbf{S}_{i1} \cdot \mathbf{S}_{i2} - t \sum_{\langle ij \rangle} |E_j\rangle \left(|00\rangle_i \langle 00|_j + \sum_m |1m\rangle_i \langle 1m|_j \right) \langle E_i|$$

Heisenberg terms Hopping term = exchange of exciton $|E\rangle$ with magnetic states $|s, m\rangle$

So what are the **dynamics** of such a bilayer exciton in a strongly correlated system? We know that if the exciton moves, the antiferromagnetic order gets frustrated.



As a result, the exciton is effectively **confined** by the antiferromagnetic spin background. As a signature, the exciton spectral function (measurable with optical absorption) obtains a second **Ising confinement** peak above the first exciton peak. [2]



What happens at any finite exciton density?

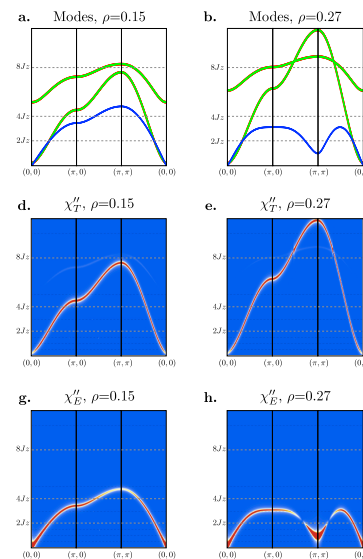
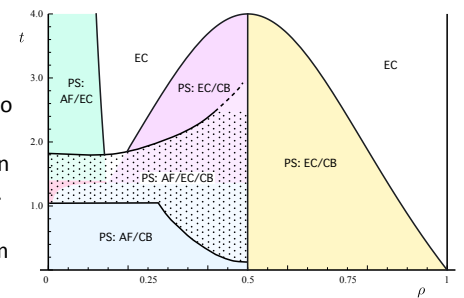
When there is a finite density of excitons, there is strong competition between three different phases:

- the **antiferromagnet (AF)** in the absence of excitons;
- the **exciton condensate (EC)**. For large exciton hopping, excitons will form a Bose condensate. The wavefunction of the condensate is given by a superposition on each rung of a spin singlet with an exciton:

$$\prod_i (u_i + v_i \hat{E}_i^{\dagger}) |00\rangle$$

- the **checkerboard phase (CB)** at half-filling, where the strong exciton-exciton repulsion leads to an exciton Wigner crystal.

However, for most parts of the phase diagram the competition leads to **macroscopic phase separation** between these three phases. (See the ground state phase diagram on the right.)



The **collective modes** (see left pictures, for $t = 2$ eV) of the exciton condensate consist of triplet excitations (middle row) and the exciton superfluid phase mode (bottom row).

The surprise that the triplet modes are characterized by a bandwidth determined by the **superfluid density of the exciton condensate**.

Therefore the superfluid density can be measured in spin fluctuation spectra (neutron scattering). [3]



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References

- [1] Eisenstein and MacDonald, *Nature* **432**, 691 (2004)
- [2] Rademaker, Wu, Hilgenkamp and Zaanen, *EPL* **97**, 27004 (2012); Rademaker, Wu and Zaanen, *New J. Phys.* **14**, 083040 (2012)
- [3] Rademaker, Hilgenkamp and Zaanen, to be published.

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