

# Quantum Thermalization and the Expansion of Atomic Clouds

Louk Rademaker and Jan Zaanen

KITP, UCSB and Leiden University, The Netherlands

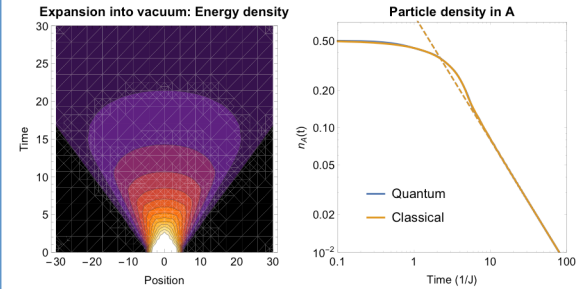
Reference: arXiv:1703.02489.

## I. Expansion of an Atomic Cloud: Classical or Quantum?

**Motivation:** In the cold atomic ‘time-of-flight’ measurements, one suddenly releases the confining potential and the atomic cloud expands. Typically, it has been assumed that this expansion is governed by a purely classical Newtonian or wave kinematics. Naively, it is not at all obvious why this works. After all, before releasing the trapping potential, one may be in a quantum regime with Bose condensation or Fermi-degeneracy. How can these atoms suddenly behave like classical canon balls?

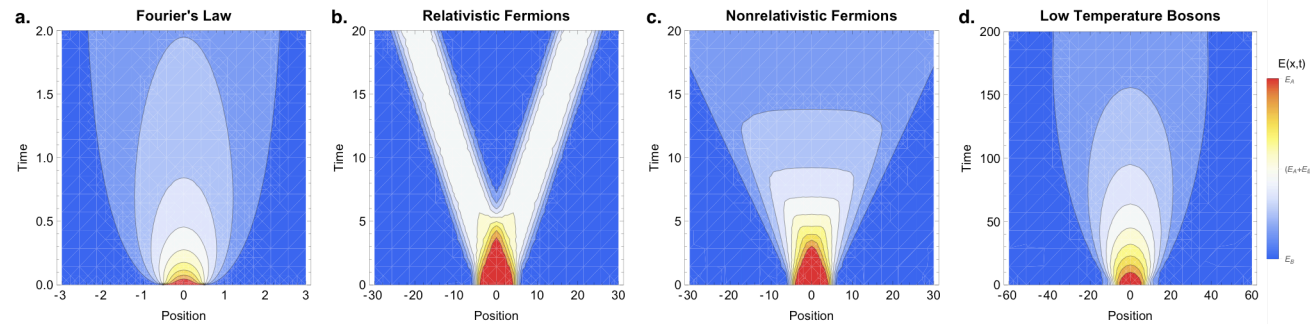
**Main question:** *Are there instances where the quantum time evolution of a macroscopic system is qualitatively different from the equivalent classical system?*

## II. Expansion into the vacuum



The expansion of an atomic cloud into vacuum is the same when computed fully quantum-mechanically or classically!

## III. Expansion into a cold bath



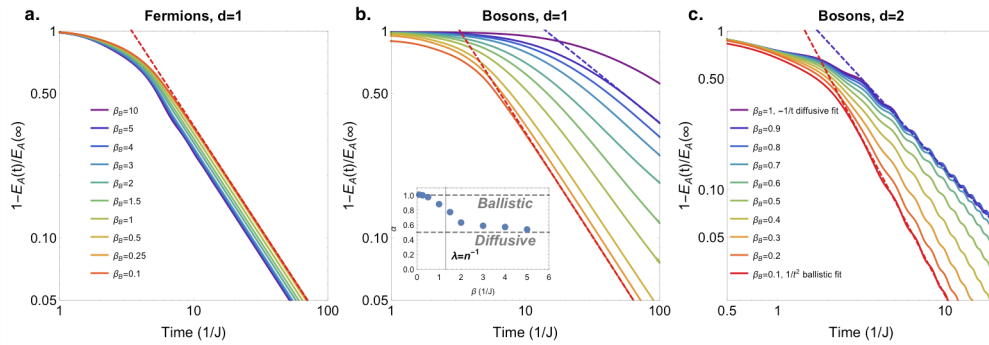
We start with a hot subsystem A in a cold bath. **How will A thermalize?**

**Classical** thermalization is diffusive,  $\partial_t T = D \nabla^2 T$  so  $\Delta T \sim t^{-d/2}$ .

**Relativistic fermions** thermalize instantaneously.

**Nonrelativistic fermions** thermalize ballistically,  $\Delta E \sim t^{-d}$ .

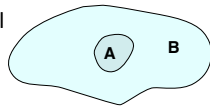
But **bosons** show a crossover from ballistic behavior, at high bath temperatures, to **diffusive at low bath temperatures!**



## IV. Method of Modular Hamiltonian

Separate a system into two parts **A** and a bath **B**. The initial density matrix is a product of a thermal state in A and a thermal state in B.

$$\rho_0 = \frac{1}{Z_A Z_B} e^{-\beta_A \mathcal{H}_A} \otimes e^{-\beta_B \mathcal{H}_B}$$



It's easier to compute the modular Hamiltonian:  $\mathcal{M} = -\log \rho$

For noninteracting particles, the time evolution is simply:

$$\mathcal{M}(t) = \sum_{kk'} m_{kk'} e^{-i(\xi_k - \xi_{k'})t} c_k^\dagger c_{k'} + \log Z$$

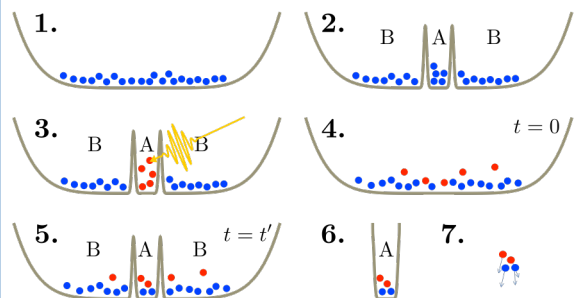
The modular matrix decays as a **ballistic** powerlaw,

$$\Delta m_{j,j+1}(t \gg 1) = 2\Delta\beta(0) \left(\frac{L_A - 1}{2\pi Jt}\right)^d \sim \frac{V_A}{t^d}$$

though the Greens function (and the energy) can have different time-dependence, since

$$\hat{G}(t) = [e^{\hat{m}(t)} - \eta]^{-1}$$

## V. Experimental realization



1. Trap an atom cloud

2. ‘Build a wall’ between A and B

3. Heat up system A

4. At t=0, remove wall

5. At later time, build wall

6. Remove all atoms in B

7. Measure kinetic energy in A using Time-Of-Flight